


TCC Week 3. The energy space - continuation

②-property: $\exists 0 < \lambda \in L^1_{loc}(\Omega), \lambda^+ \in L^\infty_{loc}(\Omega)$

$$E_V(\psi) = \int_{\Omega} |\nabla \psi|^2 + V\psi^2 \geq \int_{\Omega} \lambda(x) \psi^2 \quad \forall \psi \in C^\infty_c(\Omega)$$

Remark: $\exists u_* > 0: -\Delta u_* + V u_* \geq \underbrace{\lambda}_{\geq 0} \text{ in } \Omega \Rightarrow \lambda = \frac{\lambda}{u_*}$

Thm. If E_V satisfies ② then

$$\mathcal{D}'_V(\Omega) = {}_{ce_{\|\cdot\|_V}} C^\infty_c(\Omega), \quad \|\cdot\|_V = \sqrt{E_V(\cdot)}$$

the Hilbert space with $\langle u, v \rangle_V = \int_{\Omega} \nabla u \nabla v + \int_{\Omega} V u v$,
and $\mathcal{D}'_V(\Omega) \subset L^2(\Omega, \lambda(x) dx)$.

Remark: $(\mathcal{D}'_V(\Omega))^* \supset L^2(\Omega, \alpha^{-1}(x) dx)$ Ex.

∇ $e(\varphi) := \int_{\Omega} f \varphi$, $f \in L^2(\Omega, \alpha^{-1} dx)$ is a Bounded Functional on $\mathcal{D}'_V(\Omega)$ ∇

Thm (Lax-Milgram thm) Assume ε_V satisf. (2)

Then $\forall e \in (\mathcal{D}'_V(\Omega))^* \exists$ unique solution $u_e \in \mathcal{D}'_V(\Omega)$ such that:

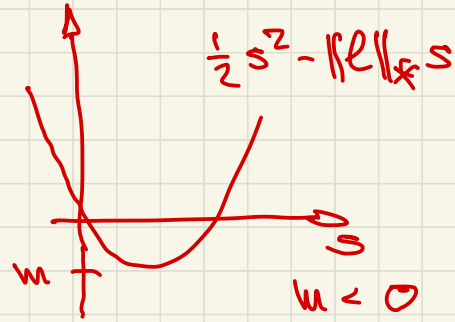
$$\langle u_e, \varphi \rangle_V = e(\varphi) \quad \forall \varphi \in \mathcal{D}'_V(\Omega).$$

Rem. if $e(\varphi) = \int_{\Omega} f \varphi$ then $-\Delta u_e + V u_e = f$ in Ω :

$$\int_{\Omega} \nabla u_e \nabla \varphi + \int_{\Omega} V u_e \varphi = \int_{\Omega} f \varphi \quad \forall \varphi \in \mathcal{D}'_V(\Omega) \supset C_0^\infty(\Omega)$$

$\blacktriangle \mathcal{E}(u) = \frac{1}{2} \|u\|_V^2 - \underbrace{e(u)}_{\langle e, u \rangle_V} = \frac{1}{2} \|u\|_V^2 - \langle e, u \rangle_V \geq$
 $\geq \frac{1}{2} \|u\|_V^2 - \|e\|_* \|u\|_V \geq m$

— $\mathcal{E}$ is Bounded below



$\mathcal{E}(u) \rightarrow +\infty$ if $\|u\|_V \rightarrow \infty$ — coercive!

Assume $(u_n) \subset \mathcal{D}'_V(\Omega)$ — minimising sequence

$$\mathcal{E}(u_n) \rightarrow m = \inf_{\mathcal{D}'_V(\Omega)} \mathcal{E}$$

1) Since $\mathcal{E}$ is coercive, $\|u_n\|_V \leq M$

2) $\|u_n\|_V \leq M \Rightarrow$ converges weakly to u_0
(up to a subsequence)

3) $\mathcal{E}(u)$ is weakly lower semicontinuous,
 $u_n \rightharpoonup u_0 \Rightarrow \mathcal{E}(u_0) \leq \liminf_{n \rightarrow \infty} \mathcal{E}(u_n)$

$$\geq m = \inf E$$

m

$$\Rightarrow \mathcal{E}(u_0) = m$$

$\Rightarrow u_0$ is a minimum of $\mathcal{E}$.

$$u) \quad \mathcal{E}(u) = \underbrace{\frac{1}{2} \|u\|_V^2}_{\text{strictly convex}} + \underbrace{e(v)}_{\text{linear} \Rightarrow \text{convex}} - \text{strictly convex}$$

$$\Rightarrow \mathcal{E}(\lambda u + (1-\lambda)v) < \lambda \mathcal{E}(u) + (1-\lambda)\mathcal{E}(v), \quad 0 < \lambda < 1.$$

Ex. $\|u\|^2$ is strictly convex.

$\Rightarrow u_0$ is unique minimiser

(Otherwise $\mathcal{E}(\lambda u_0 + (1-\lambda)v_0) < \lambda m + (1-\lambda)m = m$)

5) Since u_0 is the minimiser,

$$\underbrace{\langle u_0, \varphi \rangle = e(\varphi)}_{\text{Euler-Lagrange eq.}} \quad \forall \varphi \in \mathcal{D}'_V(\Omega)$$

1) $\mathcal{D}'_0(\mathbb{R}^N)$, $N \geq 3$ — is the energy space

$$\text{for } \int_{\mathbb{R}^N} |\nabla u|^2 \quad \left(\lambda(x) = (1+|x|^2)^{-\frac{N-2}{2}} \right)$$

$\Rightarrow -\Delta u = f$ in $\mathbb{R}^N$ has unique solution
 $\forall f \in L^2\left(\mathbb{R}^N, (1+|x|^2)^{\frac{N-2}{2}}\right) \not\subset L^2(\mathbb{R}^N)$

$$2) H^1(\mathbb{R}^N), \quad \int |\nabla u|^2 + \int u^2 \geq \int |\nabla u|^2 - \textcircled{2}!$$

$$H^1(\mathbb{R}^N) \subset \mathcal{D}'_0(\mathbb{R}^N)$$

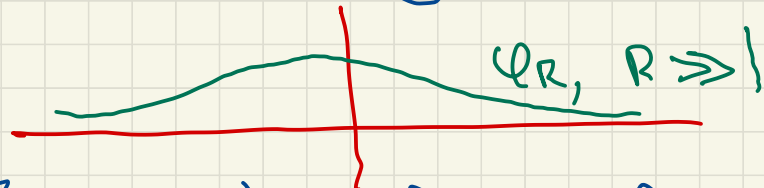
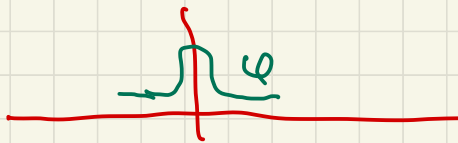
$$H^1(\mathbb{R}^N) = \mathcal{D}'_0(\mathbb{R}^N) \cap L^2(\mathbb{R}^N)$$

$$\Rightarrow (H^1(\mathbb{R}^N))^* \supset L^2(\mathbb{R}^N)$$

$$3) \int_{\mathbb{R}^N} |\nabla u|^2 - \int_{\mathbb{R}^N} u^2 \quad \text{--- } \textcircled{2} ? \quad \text{--- no!}$$

◀ Scaling argument. Take $\psi \in C_0^\infty(\mathbb{R}^N)$

$$\psi_R(x) = \psi\left(\frac{x}{R}\right) \quad \text{--- rescaling of } \psi$$



$$\int_{\mathbb{R}^N} \psi_R^2(x) dx = \int_{\mathbb{R}^N} \psi^2\left(\frac{x}{R}\right) R^N d\left(\frac{x}{R}\right) = R^N \int_{\mathbb{R}^N} \psi^2$$

$$\begin{aligned} \int_{\mathbb{R}^N} |\nabla \varphi_R(x)|^2 &= \int |\nabla \varphi(\frac{x}{R})|^2 = \int |\frac{1}{R} \nabla \varphi|^2 R^N d(\frac{x}{R}) \\ &= R^{N-2} \int |\nabla \varphi|^2 \end{aligned}$$

$$\begin{aligned} \int |\nabla \varphi_R|^2 - \int |\varphi_R|^2 &= R^{N-2} \underbrace{\int |\nabla \varphi|^2}_A - R^N \underbrace{\int \varphi^2}_B = \\ &= AR^{N-2} - BR^N \rightarrow -\infty \text{ as } R \rightarrow \infty \end{aligned}$$

$$\Rightarrow \boxed{\inf \sigma(-\Delta) = 0} \quad (\text{not needed for us})$$

Ex. Ω — bounded domain.

$$\text{Then } \int |\nabla \psi|^2 \geq \lambda_1 \int \psi^2 \quad \forall \psi \in C_0^\infty(\Omega)$$

$\lambda_1 = \lambda_1(\Omega) > 0$

$$\Rightarrow \int |\nabla u|^2 - \lambda \int u^2 \geq \underbrace{(\lambda_1 - \lambda)}_{> 0} \int u^2 \quad - \textcircled{\lambda} \quad \forall \lambda < \lambda_1$$

Remark: If $\lambda = \lambda_1 \Rightarrow \int |\nabla \psi|^2 - \lambda_1 \int \psi^2 = 0$,
 $\psi_1 > 0$ — principal eigenvalue!

According to Agmon and Pinchover

1) $-\Delta + V$ is subcritical if it satisfies (a)
(think of $-\Delta - \alpha$ on Ω with $\alpha < \alpha_1$)

2) $-\Delta + V$ is critical if (a) fails but
 $\int |\nabla \varphi|^2 \geq \int V \varphi^2 \quad \forall \varphi \in C_0^\infty(\Omega)$

(think of $-\Delta - \alpha_1$ on Ω or $-\Delta - \frac{C_H}{|x|^2}$ on $\mathbb{R}^n$)

3) $-\Delta + V$ is supercritical if $\exists \varphi \in C_c^\infty(\Omega)$
 $\int |\nabla \varphi|^2 + \int V \varphi^2 < 0.$

Example: $N \geq 3$

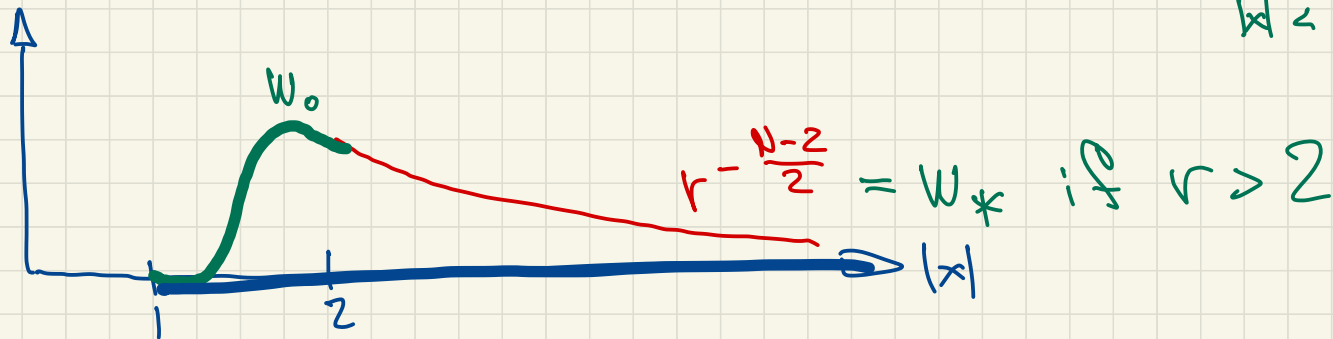
$$\int_{\mathbb{R}^N} |\psi|^2 - \int_{\mathbb{R}^N} \frac{C_H}{|x|^2} \psi^2 \geq \frac{1}{4} \int_{\mathbb{R}^N} \frac{\psi^2}{|x|^2 \log^2 |x|}$$

Consider $D'_{\frac{C_H}{|x|^2}}(\mathbb{R}^N, \overline{B}_1)$

Take $w(x) = \begin{cases} |x|^{-\frac{N-2}{2}}, & |x| > 2 \\ w_0(x), & 2 < |x| < 1 \end{cases}$

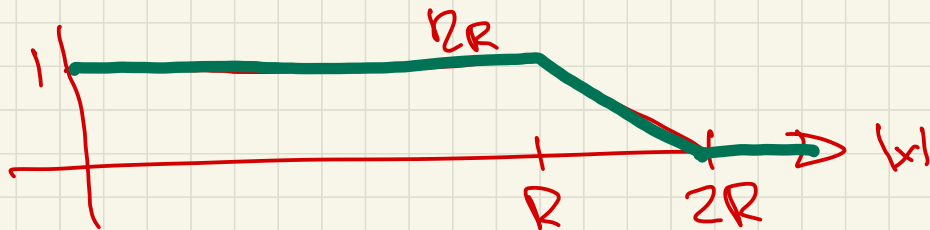
w_0 is such that w is C^2 , $w \geq 0$ and $w(x) = 0$, $|x| < \frac{1}{2}$

λ $\forall \psi \in C_0^\infty(\mathbb{R}^N \setminus \overline{B}_1)$



$$w_R = 2R w$$

comp. supp.



$$\int |w_R|^2 - \int \frac{C_H}{|x|^2} w_R^2 =$$

$$\left(-\Delta - \frac{C_H}{|x|^2} \right) r^{-\frac{N-2}{2}} = 0$$

w_*

AAP!

AAP

$$= \int \left| \Delta \frac{w_R}{w_*} \right|^2 w_*^2 =$$

$$= \underbrace{\int_0^2 \left| \Delta \frac{w_R}{w_*} \right|^2 w_*^2}_{C_1} + \underbrace{\int_2^{2R} \left| \Delta \frac{2R w_*}{w_*} \right|^2 w_*^2}_{C_2} =$$

$$\approx C_1 + \frac{C_2}{R^2} \int_R^{2R} (r^{-(N-2)}) r^{N-1} dr =$$

$$= c_1 + c_2 \frac{1}{R^2} (\log 2R - \log R) =$$

$$= c_1 + c_2 \frac{\log 2}{R^2} \xrightarrow{R \rightarrow \infty} c_1$$

$$\Rightarrow w \in \mathcal{D}'_{\frac{C_H}{|x|^2}}(\mathbb{R}^N, \overline{B}_1)$$

$$E_v(w) = \int_{\mathbb{R}^N} |dw|^2 - \int_{\mathbb{R}^N} \frac{C_H}{|x|^2} w^2 = c_1$$

But $\underbrace{\int_{-\infty+\infty} |w_R|^2}_{\rightarrow +\infty} - \underbrace{\int_{-\infty-\infty} \frac{C_H}{|x|^2} w_R^2}_{\rightarrow -\infty},$

so $\int |w|^2 - \int \frac{C_H}{|x|^2} w^2$ is not defined!

Yet $E_v(w) = C_1 < \infty$, well defined
 $w \in \mathcal{D}_{\frac{C_H}{|x|^2}}^1(\mathbb{R}^N \setminus \overline{B_1})$ but $w \notin \mathcal{D}_0^1(\mathbb{R}^N \setminus \overline{B_1})$